

# Searching for new dynamics in strongly coupled lattice gauge theory

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# Outline of talk

- Motivation: technicolor and walking technicolor
- Phase diagram non-supersymmetric gauge theory
- Lattice  $SU(2)$  with 2 flavors of adjoint quark
- Simulation results
- Conclusions and outlook

# Troubles with Higgs ...

Standard model Higgs has many problems:

- Quadratically unstable against radiative corrections.
- Triviality
- Symmetry breaking “put in by hand”. What sets scale ?

One attractive solution: dispense with elementary Higgs, invoke new **strong dynamics** to break EW symmetry:

**Technicolor**

# Basic idea

- Assume new strong dynamics at scale  $\Lambda_{rmTC} = O(1)$  TeV causes condensation of new techni-quarks
- If electroweak symmetry embedded into chiral-flavor symmetries of techniquarks will precipitate EW breaking
- Massive EW bosons produced from would-be Goldstone's associated with broken chiral symmetries

# Advantages

- Follows example of QCD: asymptotically free gauge theories naturally develop a scale  $\Lambda \ll M_{\text{GUT/Planck}}$  with chiral symmetry breaking.
- No quadratic fine tuning. No triviality problem.

## Problems:

- Constraints: How many techni-quarks, how many colors, what representation ?
- Phenomena arise from strong dynamics: how to make predictions ?
- What about masses of quarks/leptons – need to couple techni-quarks to usual Standard Model particles – extended technicolor.
- Understand flavor hierarchy in terms of dynamics at

$\Lambda_{\text{FTC}} ?$

# Why QCD won't do

- Scaled up versions of QCD at odds with precision EW observables.
- In generic ETC theories – effective 4-quark operators yield FCNC – suppress by large  $\Lambda_{\text{ETC}}$ . But then fermion masses too small !

Solution: dynamics must be **very** different from QCD

One possibility:  $g_{\text{TC}}$  must **walk** in  $\Lambda_{\text{TC}} \rightarrow \Lambda_{\text{ETC}}$

# How can this arise ?

- Slow running requires theory lie near zero of the  $\beta$ -function. Theory near conformal.
- Can we find realistic walking gauge theories ?
- More generally: what types of dynamics are possible in 4d gauge theory ...?

# Phases of gauge theory

Distinguish using long distance potential between static heavy sources.

- Coulomb  $V(r) = \frac{1}{r}$ . Conformally invariant.
- Free electric  $V(r) = \frac{1}{r \ln r}$ . Eg QED.
- Free magnetic  $V(r) = \frac{\ln r}{r}$
- Higgs  $V(r) = \text{const}$
- Confining  $V(r) = \sigma r$

All realized in supersymmetric GT. Non-supersymmetric ?  
Can we find a region in  $(N_c, N_f, R)$  where theory develops I.R conformal fixed point ?

# Phase diagram

Variety of approx analytical approaches used

Conformal window bounded below by chiral symmetry breaking and above by loss of asymptotic freedom.

# Minimal theories

Notice: Using symmetric/adj reps significantly lowers the number of flavors needed for walking Minimal theory  $SU(2)$  with 2 adjoint flavors is (near)conformal!

- Ladder:  $N_f \sim, 2.1$
- Conjectured all order beta function – theory **already** in conformal window (Sannino)

Minimal walking technicolor MWT possible (Sannino).

# Lattice approach - glue

Use lattice simulation to look for conformal phase.  
Employ standard Wilson gauge action

$$S_G = -\frac{\beta}{2} \sum_x \sum_{\mu > \nu} \text{ReTr} \left( U_\mu(x) U_\nu(x + \hat{\mu}) U_\mu^\dagger(x + \hat{\nu}) U_\nu^\dagger(x) \right) ,$$

where  $U_\mu(x) = e^{A_\mu(x)a}$  is in fundamental of  $SU(2)$  and

$$\beta = \frac{4}{g^2}$$

Show  $S_G \rightarrow \int d^4x F_{\mu\nu}^2 + O(a)$

# Lattice approach - quarks

Several choices. Use Wilson. Fast, no rooting problem.

$$S_F = -\frac{1}{2} \sum_x \sum_\mu \bar{\psi}(x) \left( V_\mu(x) (I - \gamma_\mu) \psi(x + \mu) + V_\mu^T(x - \mu) (I + \gamma_\mu) \psi(x - \mu) \right) + \sum_x (m + 4) \bar{\psi}(x) \psi(x) .$$

with adjoint links

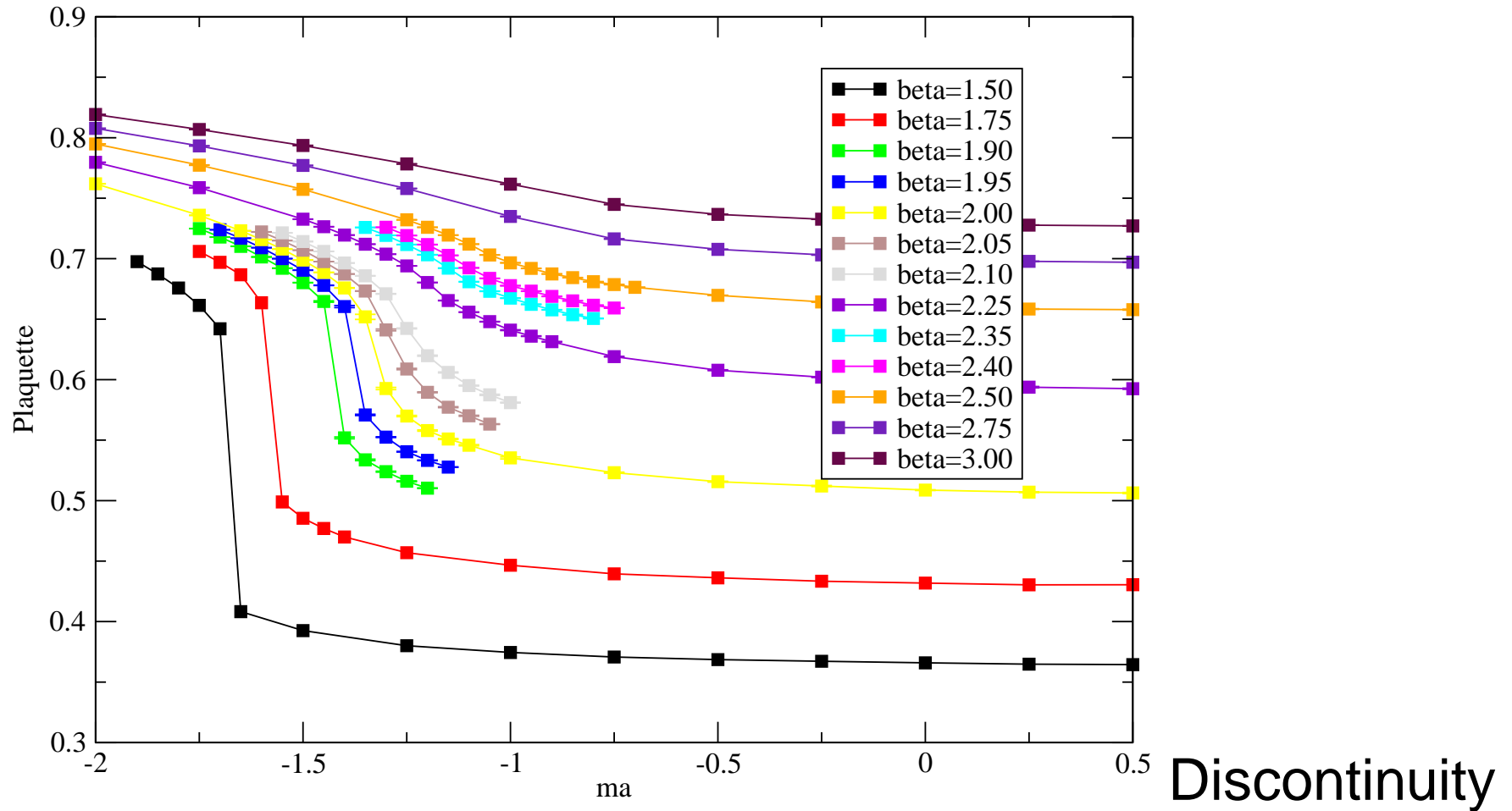
$$V_\mu^{ab}(x) = \text{Tr} \left( {}^a U_\mu(x) {}^b U_\mu^\dagger(x) \right) ,$$

Drawbacks: need to tune bare quark mass  $m$  with  $\beta$  to restore chiral symmetry.

# Numerical details

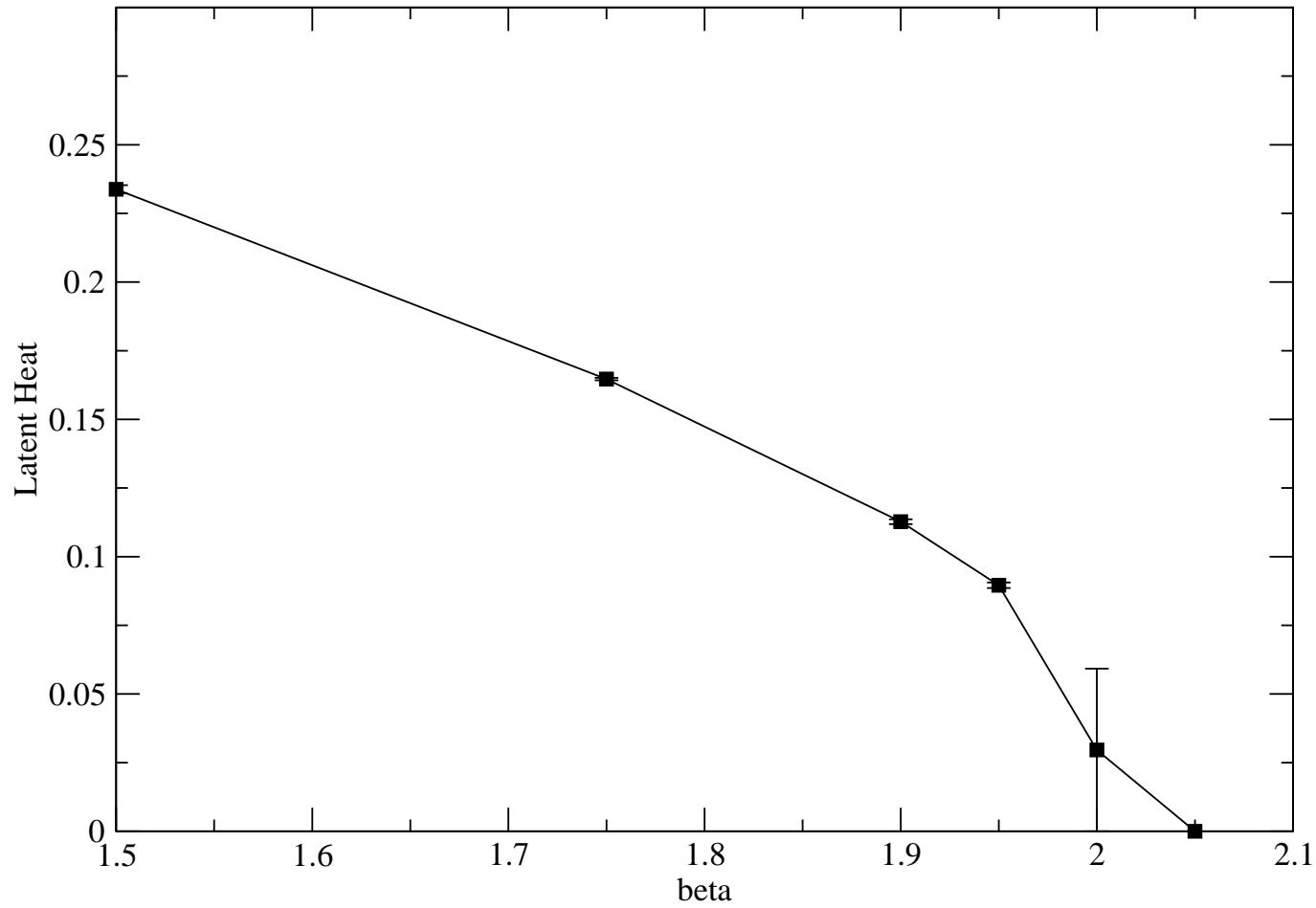
- Use HMC algorithm – efficient and exact way to handle dynamical fermions
- Implemented using a hack of CPS to allow for adjoints.
- Runs correspond to 4+ months on 128 node rack of IBM BlueGene L at RPI.
- Lattice sizes:  $8^3 \times 16$  and  $16^3 \times 32$ .
- Initial work: find phase structure in  $(\beta, m)$  plane.
- Continuous phase transitions imply a massless state and allow a continuum limit. If chiral symmetry broken – massless state is pion.
- Scan in  $ma$  at fixed  $\beta$  for minimum pion(rho) mass and critical behavior of gluon action.

# Mean gluon action



at small  $\beta$ : 1st order transition.

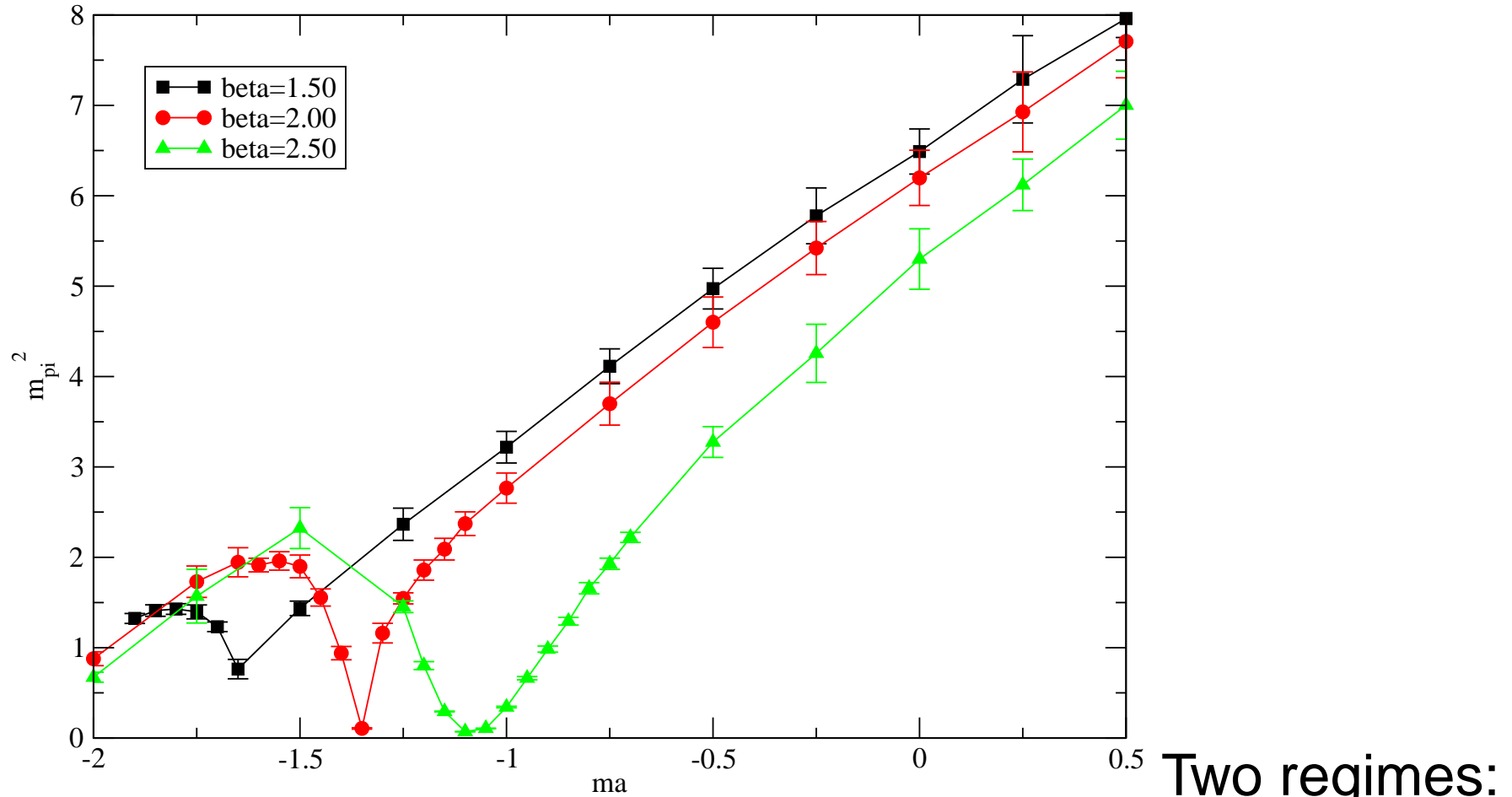
# Latent Heat



Critical end

point  $\beta = \beta_c \sim 2.0$  ?

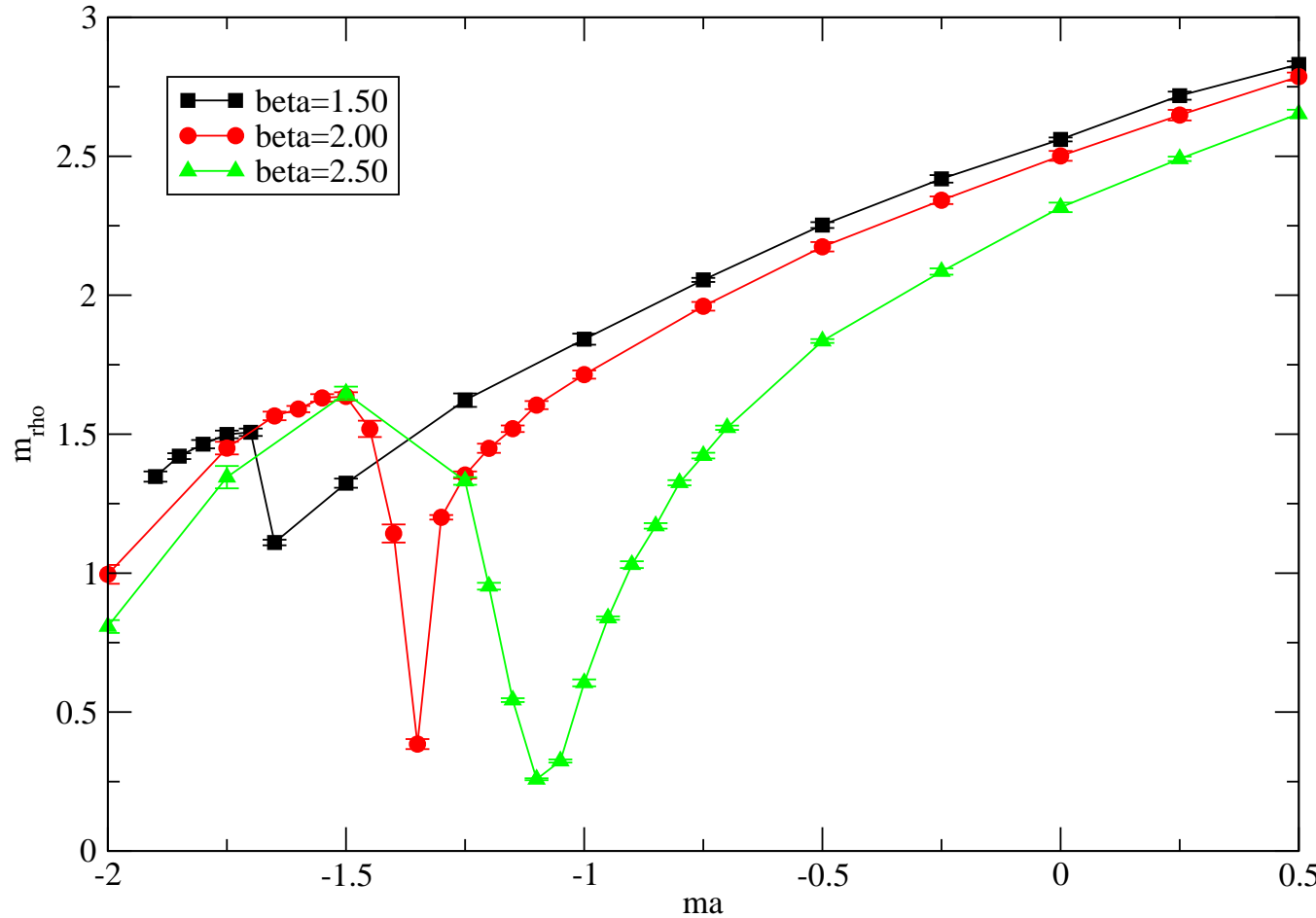
# Pion mass squared



$\beta < \beta_c$ : Goldstone behavior. Coincides with jump in gluon action  
action

$\beta > \beta_c$ : Restoration of chiral symmetry (?). Gluon action  
smooth.

# Rho mass

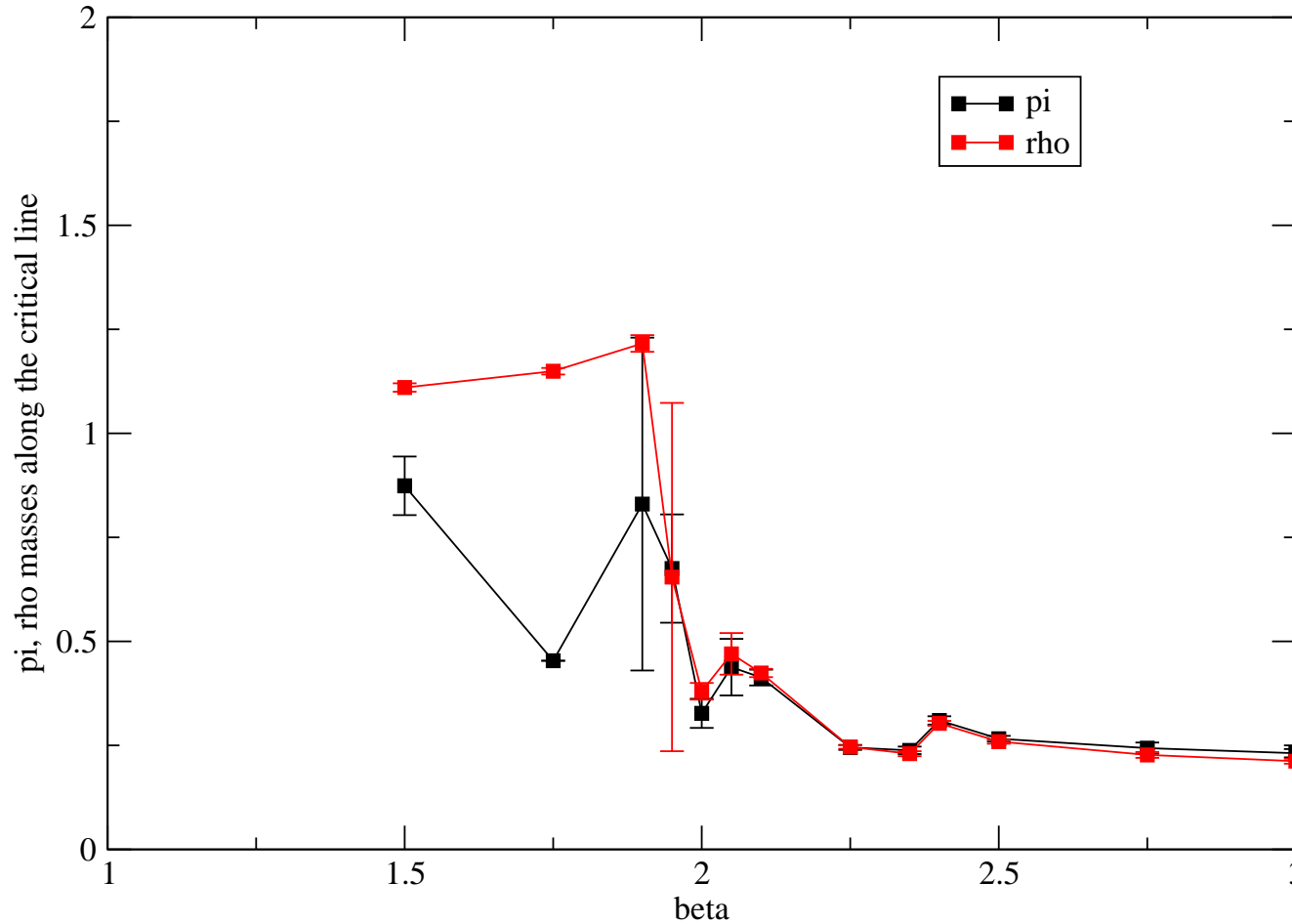


Two regimes:

$\beta < \beta_c$ : Large mass. Coincides with jump in gluon action

$\beta > \beta_c$ :  $m_\rho$  small and **independent** of bare coupling. Gluon action smooth.

# Masses along “critical” line

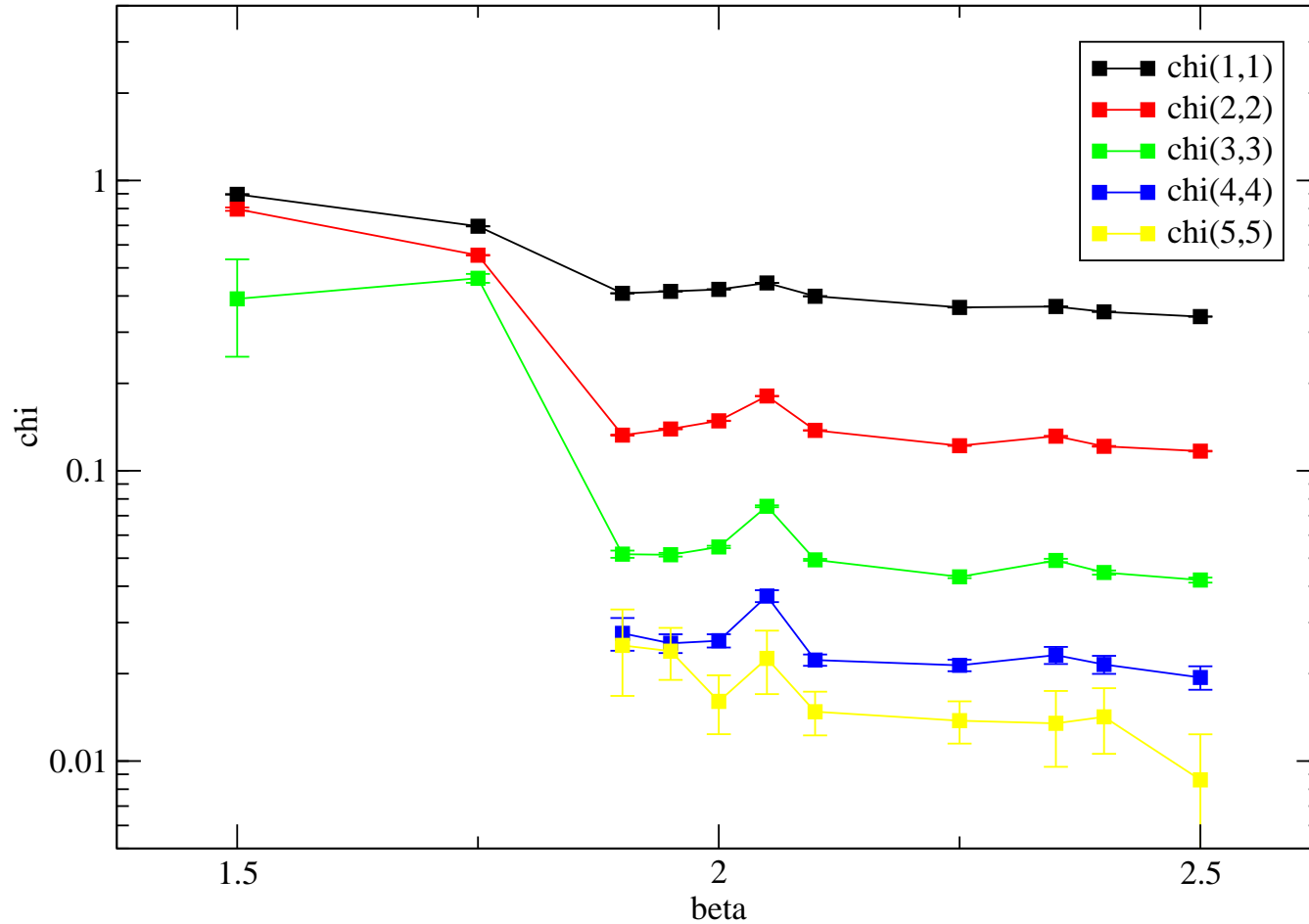


Notice:

masses degenerate for  $\beta > \beta_c$

Notice:  $\beta > \beta_c$   $m_\pi = 0.2 \sim \frac{\pi}{T} \sim$  thermal mass.

# String tension



Estimate from

Creutz ratio  $\chi(R, R)$  (ratios of Wilson loops)

If  $\frac{\chi(R, R)}{\chi(R+1, R+1)} \rightarrow 1$   $\sigma a^2 = \chi(R, R)$

Two regimes:

# Conclusions

Along line where  $m_q^r \sim 0$  see two phases.

- Confining, chirally broken phase with strong lattice artifacts  $m_\rho a, \sigma a^2 \sim O(1)$ .
- Chirally symmetric phase with  $m_\pi \sim m_\rho \sim m_{\text{thermal}}$  and  $\sigma a^2 \leq 0.0015$ . Deconfined ?
- Latter phase meson masses **independent** of bare coupling.

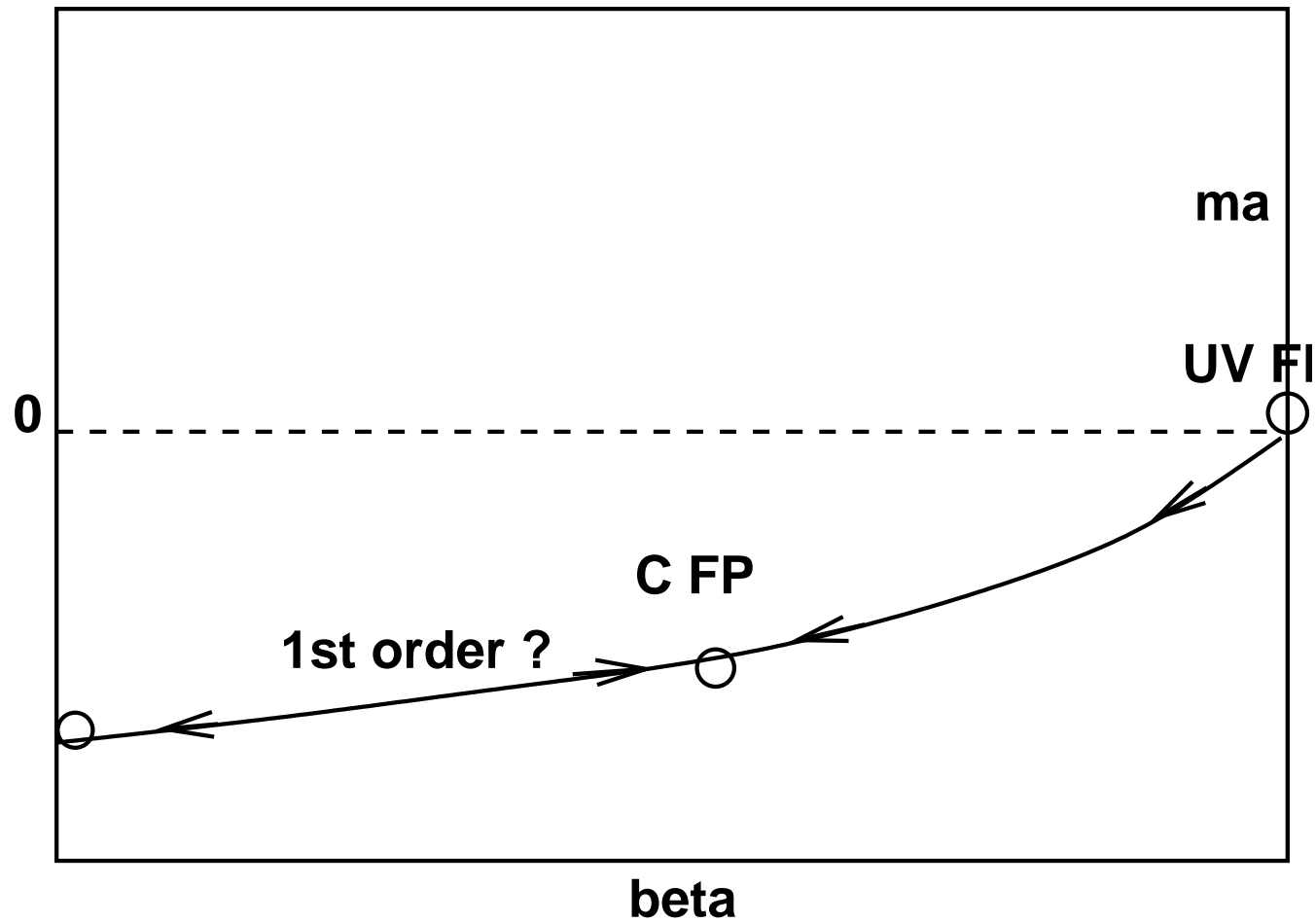
Observations consistent with existence of infrared stable fixed point

# Caveats

- Continuum limit requires careful tuning of  $\beta$  with  $a$  taking care to avoid finite size effects
- If running  $\beta(a)$  slow expect extreme sensitivity in inverse  $a = a(\beta)$ . Modest increases in  $\beta$  correspond to huge **decreases** in lattice spacing and for fixed lattice size large potential finite size effects.
- Eg. Physical box size so small system deconfines and looks quasi-free. Consistent with our data.
- Alternatively dynamical scale  $\Lambda a$  will be small and require large lattices to see asymptotic confining behavior.
- Need large lattices to resolve between these scenarios

# Possible scenario

Simplest picture for new CFP. Assume



# Conclusions/Outlook

- Have performed high statistics, fine resolution survey of phase diagram of candidate (near)conformal gauge theory –  $SU(2)$  with  $N_f = 2$  adj quarks.
- Standard LGT methods used – HMC Wilson.
- Find theory exhibits either a very slow running of the gauge coupling **or**
- Contains an IR attractive conformal fixed point.
- Try to resolve using larger lattices and step scaling techniques (current)
- Either way. Theory looks very different from QCD and shows the growing utility of lattice GT techniques in exploring non-perturbative models of BSM physics.